

Disclaimer: this is not a lecture but aim to revise some of the key-aspects:

- Luminosity, energy and their limits
 - Concepts
 - Examples

- Beam Cooling
 - Concepts
 - Examples

- Q&A and advanced lumi topics
 - Crossing angle & 'Hour Glass' effect
 - Magic of integrated luminosity

Accelerator Collider Optimisation Considerations

Where every high-energy-physics particle quest starts:



Event Rate → the frequency a given particle is created per second
Physics detectors

cross-section → probability that a given particle is created
Mother Nature defines that for us and typically depends on energy

$$\dot{N}_{event} = L \cdot \sigma_{physics}$$

Luminosity → the frequency of how often the particles are brought in to collisions
Accelerator design and operation

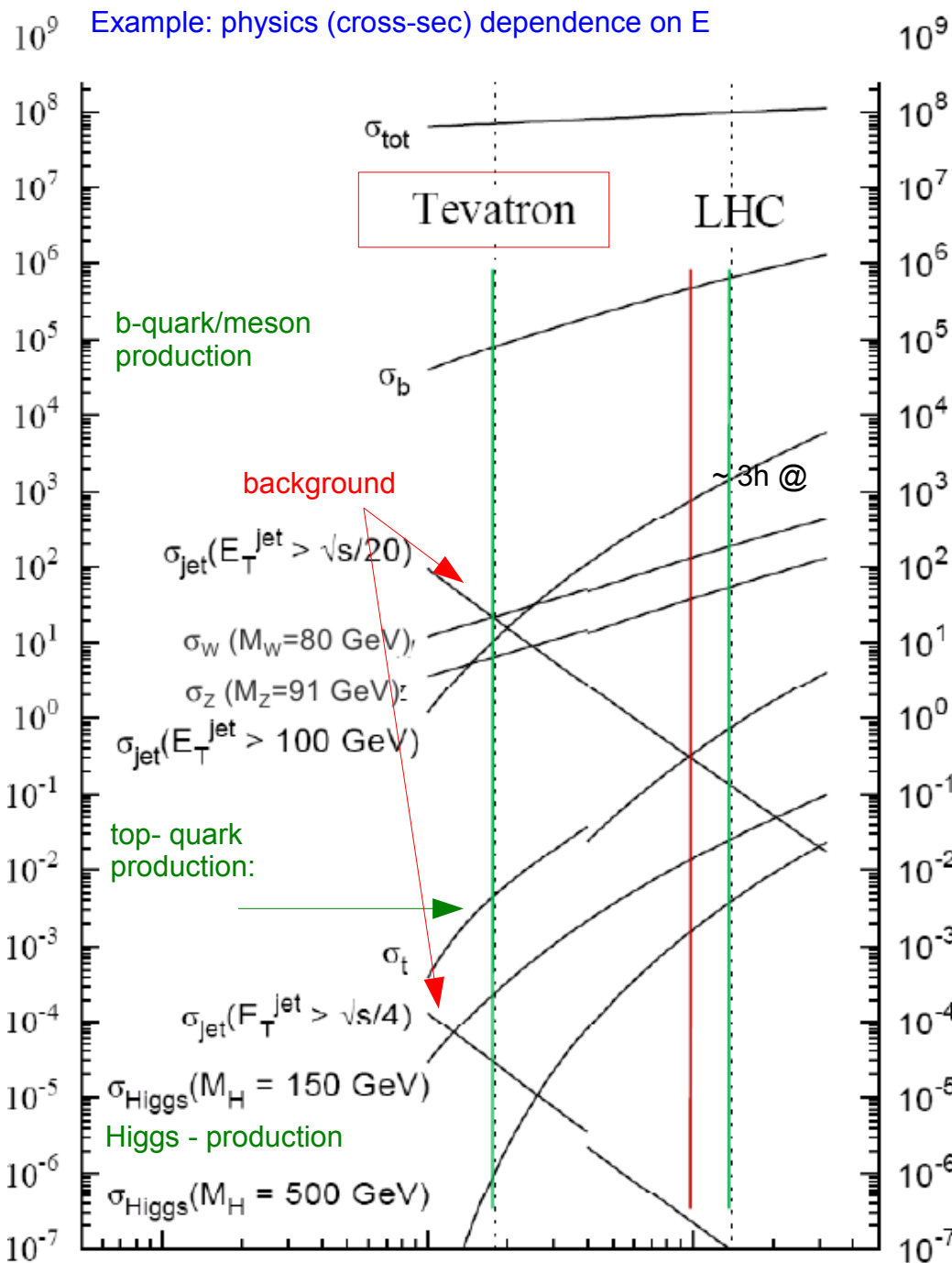
- **Push maximum peak luminosity**
 - essentially: increase number of particles inside the machine and squeeze them to a confined space to increase the probability of a collision
- **Push achievable energy E:**
 - Minimise synchrotron radiation losses: $e^+e^- \rightarrow$ hadrons collider (p^+p^\pm, \dots)
 - Choice: linear vs. circular
 - Optimise RF cavities + normal conducting magnets (CLIC,ILC)
 - Standard RF cavities + superconducting magnets (TeV,RHIC,LHC)

Why maximum Multi-TeV Collision Energies?

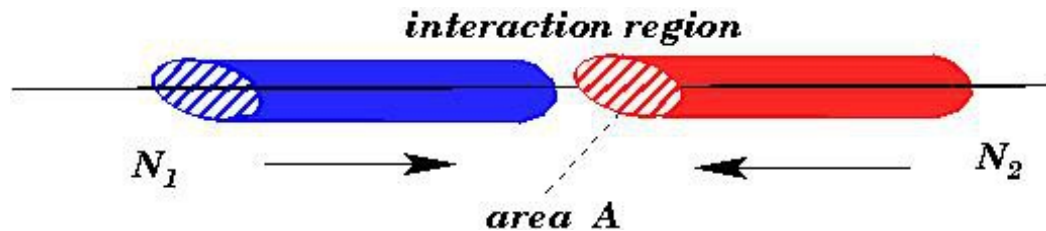
- Physicists want as high as possible
 - useful (competitive) HEP

Need to find a balance between:

- Dipole field which can be reached
- Risks associated with operating at that field
 - splices stability, thermal runaway, splice detection
- Operational efficiency of other systems
 - e.g. cryo recovery time: 5TeV vs. >6h @ 7TeV)



- Collider design:



$$L = \frac{N^2 k f_{\text{rev}}}{A} \cdot F_{\text{corr.}} = \frac{N^2 k f_{\text{rev}}}{4\pi \sigma_x \sigma_y} \cdot F_{\text{corr.}} = \frac{N^2 k f_{\text{rev}}}{4\pi \beta^* \epsilon} \cdot F_{\text{corr.}}$$

- N : number of particles per bunch,
- k : total number of bunches,
- σ_x, σ_y : hor./vert. r.m.s. beam size in IR,
- f_{rev} : revolution or repetition frequency,
- $F_{\text{corr.}}$: numerical correction factors (hour-glass, crossing angle, ...),

$$\sigma_{x,y} = \sqrt{\epsilon \beta(s) + D(s) \cdot \frac{\Delta p}{p}}$$

$$I_{\text{stored}} = Nk$$

- Warning: hadron accelerators tend to express emittance as normalised emittances!!**

- Facilitates monitoring of blow-up through the injector chain
- N.B. e-machines: emittance given by energy, damping, lattice rather than injectors

$$\epsilon^* = \frac{\epsilon}{\gamma}$$



Exercise I

– The present, from a circular Accelerator far far away...

- 1 Calculate the luminosity for a beam of 10^{11} protons per bunch circulating at 7 TeV in a ring of $C=26.658$ km circumference, assuming $\varepsilon(\beta\gamma) = 3.1 \mu\text{m rad}$ and $\beta^* = 0.55$ m. Emittance is defined for 1σ and there are 2808 bunches per beam. (assume $A=4\pi\sigma^2$)
- 2 Calculate the beam–beam tune shift per crossing for the above.
- 3 Assuming that ΔQ_{bb} should be $< 8.3 \cdot 10^{-4}$, adjust the number of bunches and intensity to maximize the luminosity within this limit. Assume $N f_b$ remains constant.



Exercise I – Solution

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■ What is needed:

- Proton mass at rest: $m_p \approx 2000 m_e \approx 0.938 \text{ GeV}/c^2$
- At 7 TeV protons are ultra-relativistic $\rightarrow \beta \approx 1 \rightarrow f = c/C$ ($c \approx 3 \cdot 10^5 \text{ km/s}$)
- Beam-beam tune shift (N.B. $r_0 \approx 0.842 \text{ fm}$):

$$\Delta Q_{bb} \approx \frac{r_0 \cdot \beta^* \cdot N}{4\pi\gamma\sigma^2} = \frac{r_0}{4\pi\gamma} \cdot \frac{N}{\epsilon} = \frac{r_0}{4\pi} \cdot \frac{N}{\epsilon^*}$$

Independent of β^* and energy!!

- Not fully “true” for LHC...

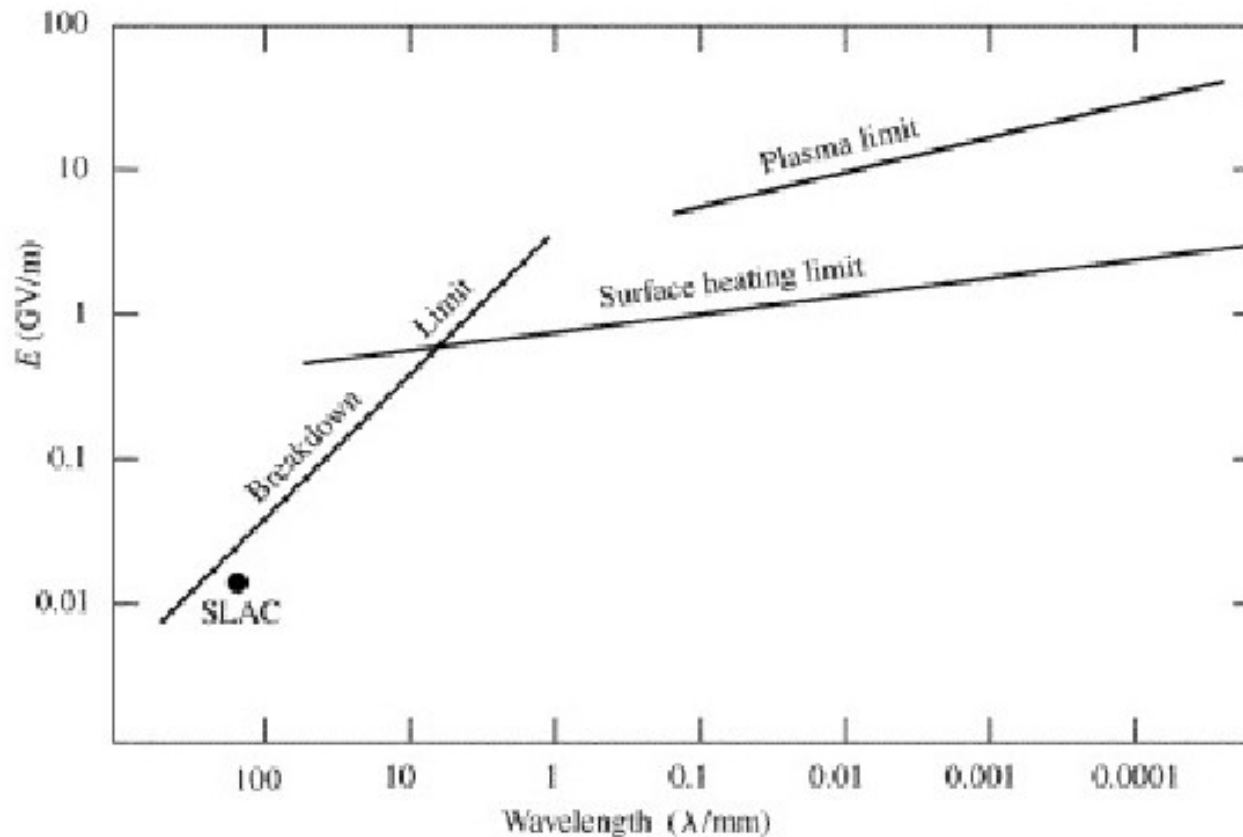
■ Lumi – design (attention 'cm' vs. ' μm ' definition):

$$L = \frac{N^2 k f_{rev}}{4\pi\sigma_x\sigma_y} \cdot F_{corr.}$$

Exercise II

– The future If you would design of the next Linear Collider...

- A linear collider seeks to achieve a centre-of-mass energy of 1000 GeV. Plot a curve of length versus field gradient and use the figure below to decide the frequency which fits a site of 25 km extent (assume a filling factor of 70%).
- Assuming a repetition frequency of 200 Hz and a mean beam radius ($\sigma_x = \sigma_y = 60$ nm), what beam intensity does the linear collider require to reach luminosity of $10^{34} \text{cm}^{-2} \text{s}^{-1}$?
- What is the average beam power?



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- Gradient = $2 \cdot 500 \text{ GV}/25 \text{ km} > 0.04 \text{ GV/m}$

- $f_{\text{rf}} = c/\lambda$, e.g. 30 cm $\rightarrow f_{\text{rf}} \approx 1 \text{ GHz}$

- Lumi-design
$$L = \frac{N^2 k f_{\text{rep}}}{4\pi \sigma_x \sigma_y} \cdot F_{\text{corr.}}$$

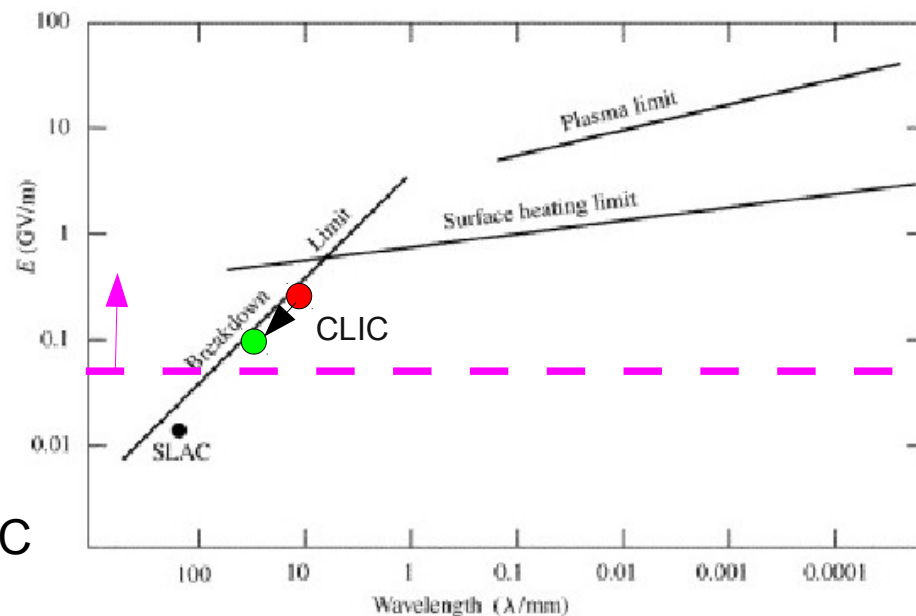
- $k=1 \rightarrow N \approx 1.5 \cdot 10^{11}$ charges = 24 nC

- N.B. 1 C(oulomb) = 1As

- Electron/Proton charge: $1.602 \cdot 10^{-19} \text{ C}$

- Power $P = N \cdot E \cdot f_{\text{rep}}$,

- here: $P = 1.5 \cdot 10^{11} \cdot 1000 \cdot 10^9 \cdot 1.602 \cdot 10^{-19} \text{ VAs} \cdot 200 \text{ Hz} \approx 5 \text{ MW}$ (per bunch)





Exercise III

– The far future...

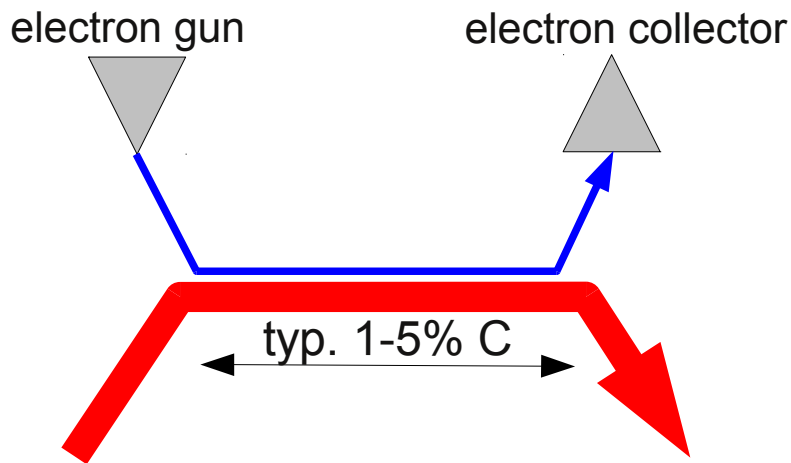
- Calculate the lifetime of a muon
 - at 50 GeV circulating in a storage ring and
 - at 4 TeV.
- Calculate the leading parameters (bending radius Q , number of periods, etc.) for a 6 T superconducting ring to store the muons.
- What would be a reasonable repetition (filling) rate for a muon collider at 4 TeV, assuming that a beam should be renewed when it has decayed by one exponential lifetime?

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-
- What is needed:
 - muon mass and life-time at rest: $m_\mu \approx 200 m_e \approx 105.7 \text{ MeV}/c^2$, $\tau_\mu = 2.2 \mu\text{s}$
 - A bit of special relativity $E = \gamma mc^2$ and time dilation $\Delta\tau' = \gamma \cdot \Delta\tau$
 - Bending radius of a dipole magnet: $p/e = (B\rho)$
 - From lattice design:
 - simple estimate $Q \approx \bar{R}/\bar{\beta}$ & $\beta_{\min, \max} = \frac{(1 \pm \sin(\mu/2))}{\sin \mu} \cdot L_{\text{cell}}$
 - e.g. (hadrons) $\Delta\mu = 90^\circ \rightarrow \bar{\beta} \approx L_{\text{cell}}$
-

- Electrons → comes easily/naturally with synchrotron radiation $\sigma_{x,y} = \sqrt{\epsilon \beta(s) + D(s) \cdot \frac{\Delta p}{p}}$
- Protons → can be created in abundance → select those fitting to small ϵ
- Secondary beams or rare isotopes, difficult to produce large numbers, e.g.: $\sim 10^6$ protons on target to create one (!) anti-proton

Electron Cooling:



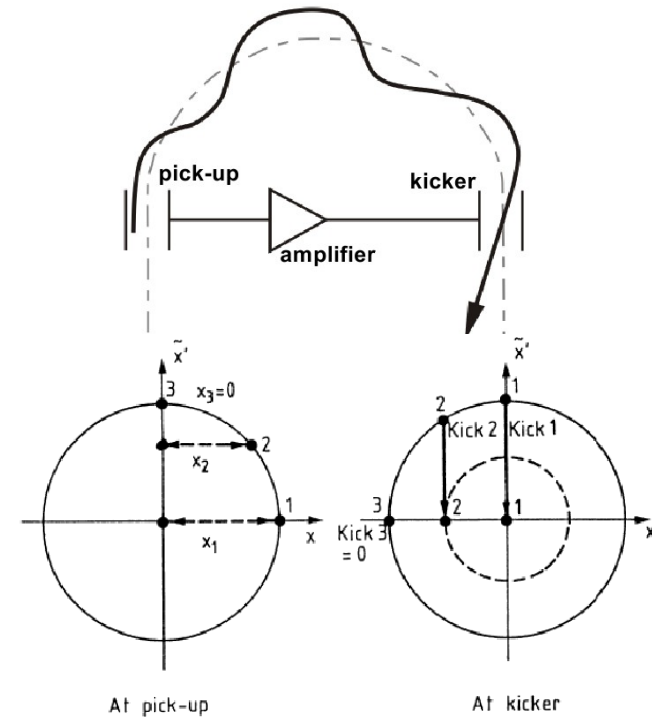
“hot” ion beam

- watch out for: alignment, RF ripple, e-space-charge, recombination

$$\tau_z \propto \frac{A}{Q^2} \frac{1}{n_e \eta} \beta^3 \gamma^5 \theta_z^3 \begin{cases} \theta_{\perp}(s) = \sqrt{\frac{\epsilon}{\beta_{\perp}(s)}} \\ \theta_{\parallel} = \beta \frac{\Delta p}{p} \end{cases}$$

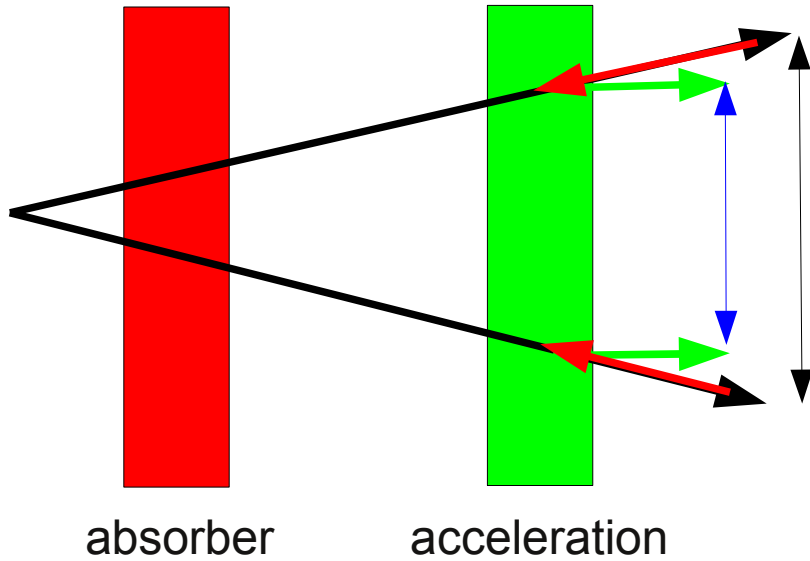
Stochastic Cooling

Van der Meer & Rubbia → Nobel Prize in '84!



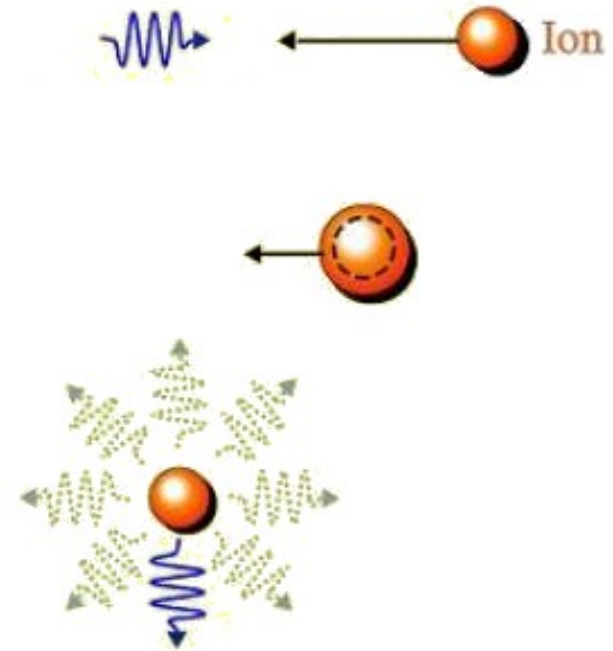
$$\frac{1}{\tau} = \frac{W}{N} \cdot \left(\underline{2g} - \underline{g^2(1+\rho)} \right)$$

- Muon cooling:



- Laser cooling – uses 'Doppler' shift ...yet another Nobel Prize in 1997

$$\omega' = \omega \gamma (1 - \beta \cos(\theta))$$



- Issue: doesn't work with too many ions...

$$\begin{aligned} \frac{d\epsilon_N}{ds} &= -\frac{1}{\beta^2 E} \frac{dE}{ds} \epsilon_N + \frac{\beta \gamma \beta_{\perp}}{2} \frac{\langle \theta_{rms}^2 \rangle}{ds} \\ &= -\frac{1}{\beta^2 E} \frac{dE}{ds} \epsilon_N + \frac{\beta_{\perp} E_s^2}{2\beta^3 m_{\mu} c^2 L_R E} \end{aligned}$$

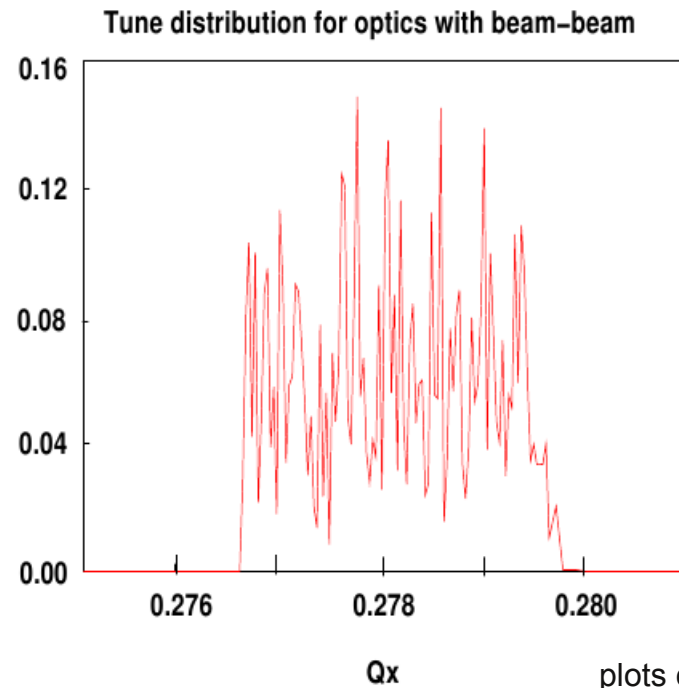
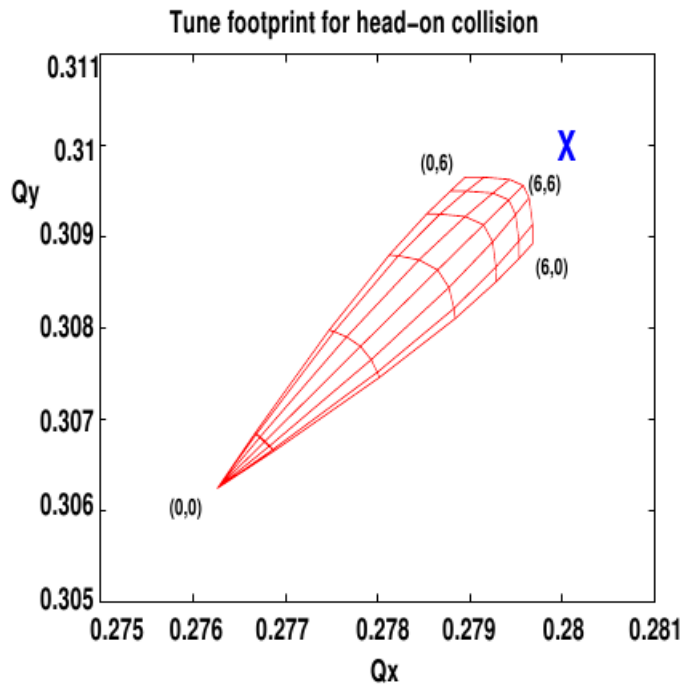
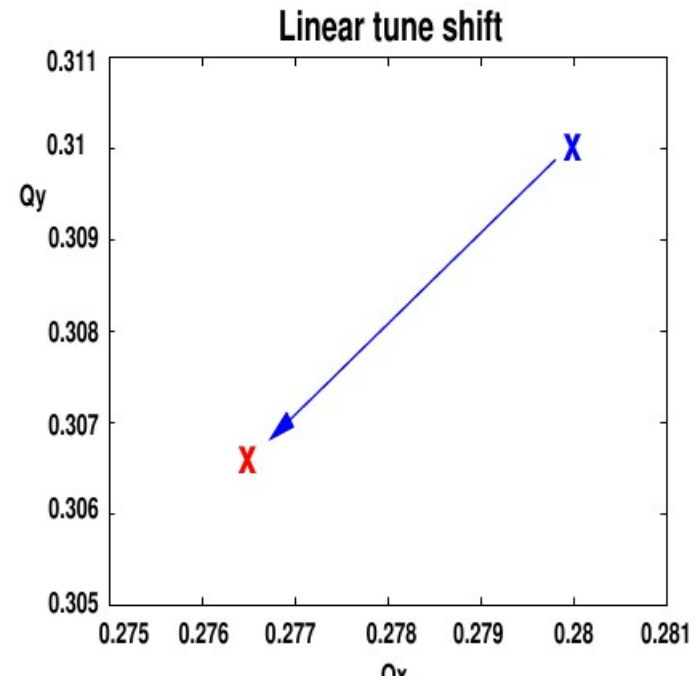
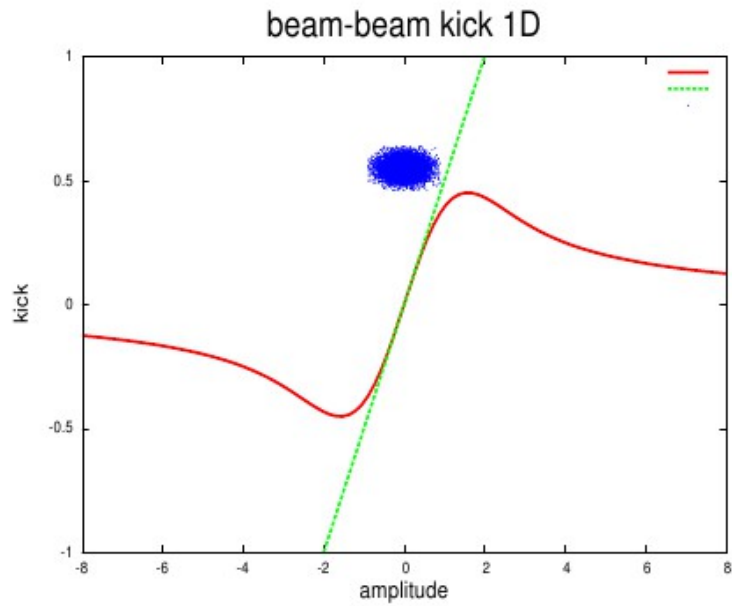
- Use the following expression for the cooling rate, and obtain an optimum value for the gain g in the presence of a signal-to-noise ratio ρ . Write down the expression for the cooling rate for this optimum gain.

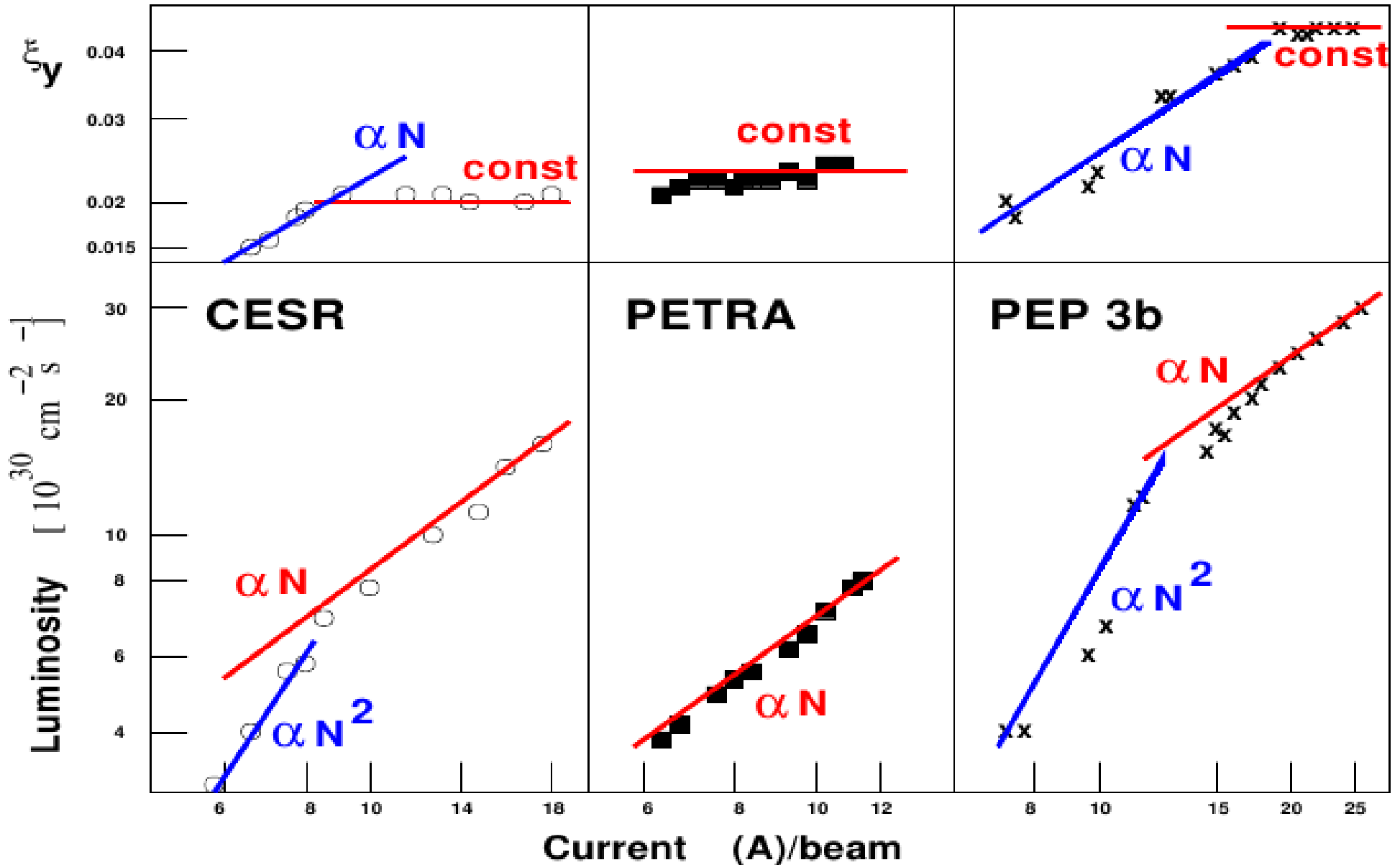
$$\frac{1}{\tau} = \frac{W}{N} \cdot (2g - g^2(1 + \rho))$$

- A cooling system is designed with a central frequency of 300 MHz. What is the sample time and how many particles will be in any one sample of 10^6 a beam of particles circulating in a ring of 25 m radius? (Assume $\beta = 0.96$)
- An electron gun has a source potential of 60 kV. Calculate the momentum of protons with the same velocity
- We wish to cool a beam of protons 5 cm in diameter to an emittance of 40π mm mrad. What is the acceptable alignment tolerance on the electron beam?
- If the proton beam has a transverse emittance of 2π mm mrad. What would be its transverse velocity at a β of 10 m? What does this represent in terms of temperature?
- The transitional state of a 100 keV Li^+ beam is excited by a laser frequency of 5485 \AA . What is the energy level difference which is excited? ($1\text{ \AA} = 0.1 \text{ nm}$, $h = 4.136 \cdot 10^{-15} \text{ eVs}$)
- Write a short computer program to simulate stochastic cooling with and without mixing.



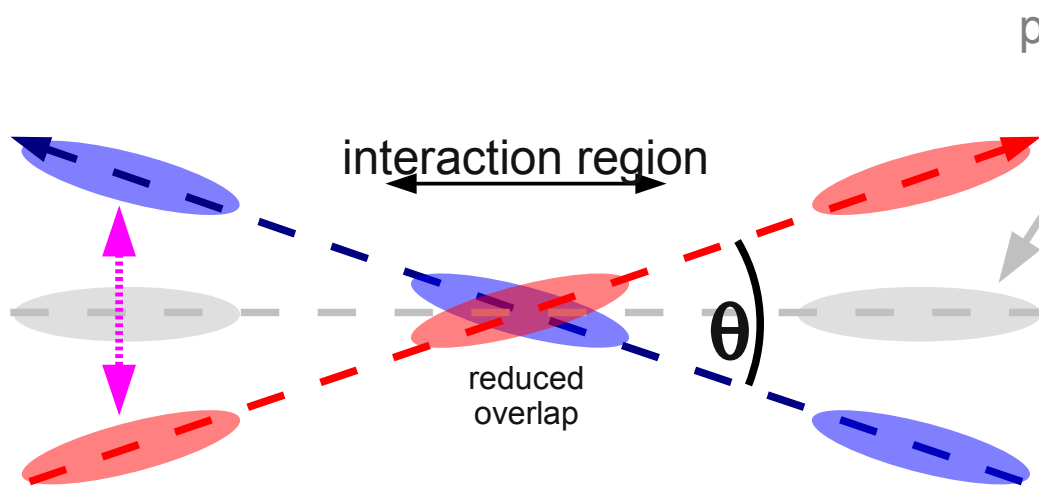
additional slides follow





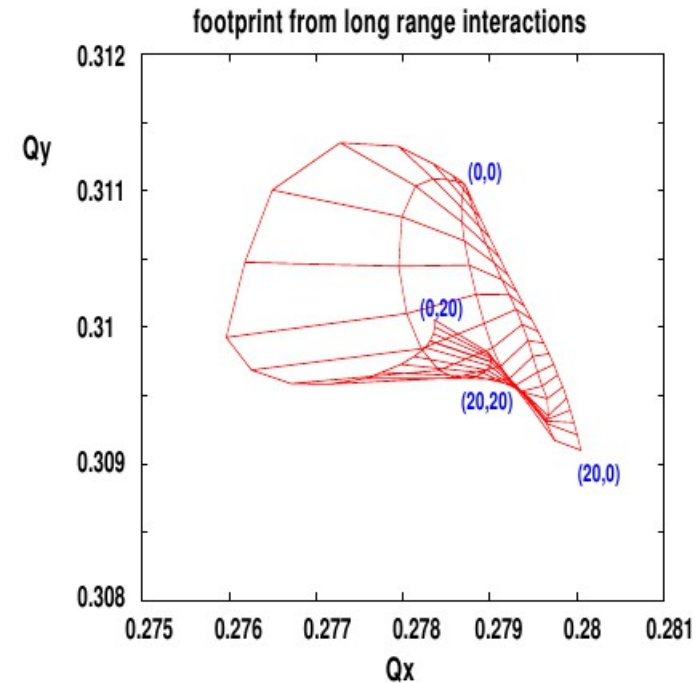
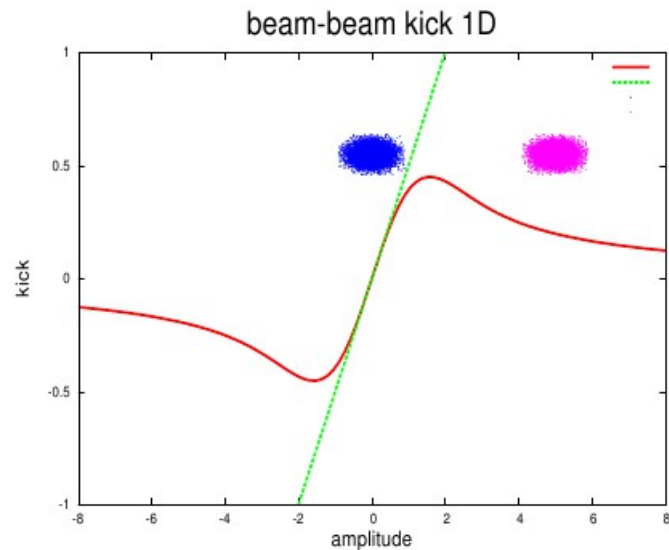
Beam-Beam in a Nutshell – Long Range

- Crossing angle θ to avoid additional parasitic crossings \rightarrow reduced bunch overlap
- “crab cavities” to compensate this effect: rotate the bunches before and after the IR

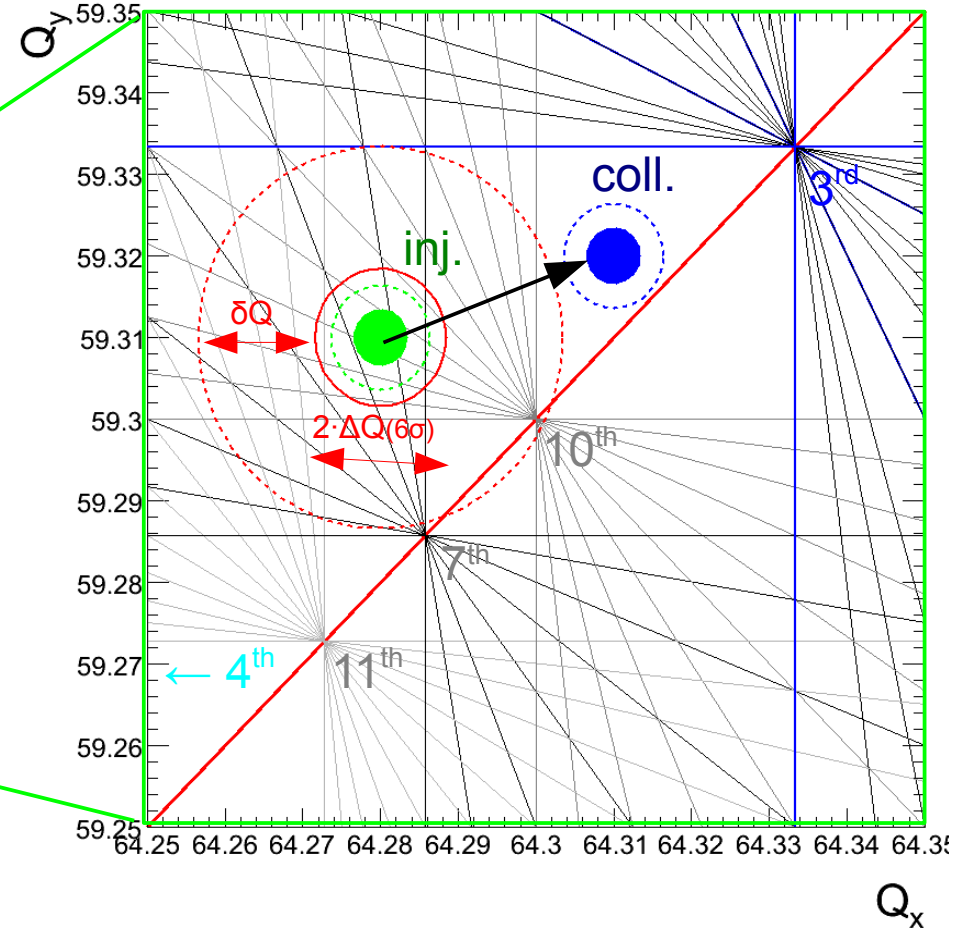
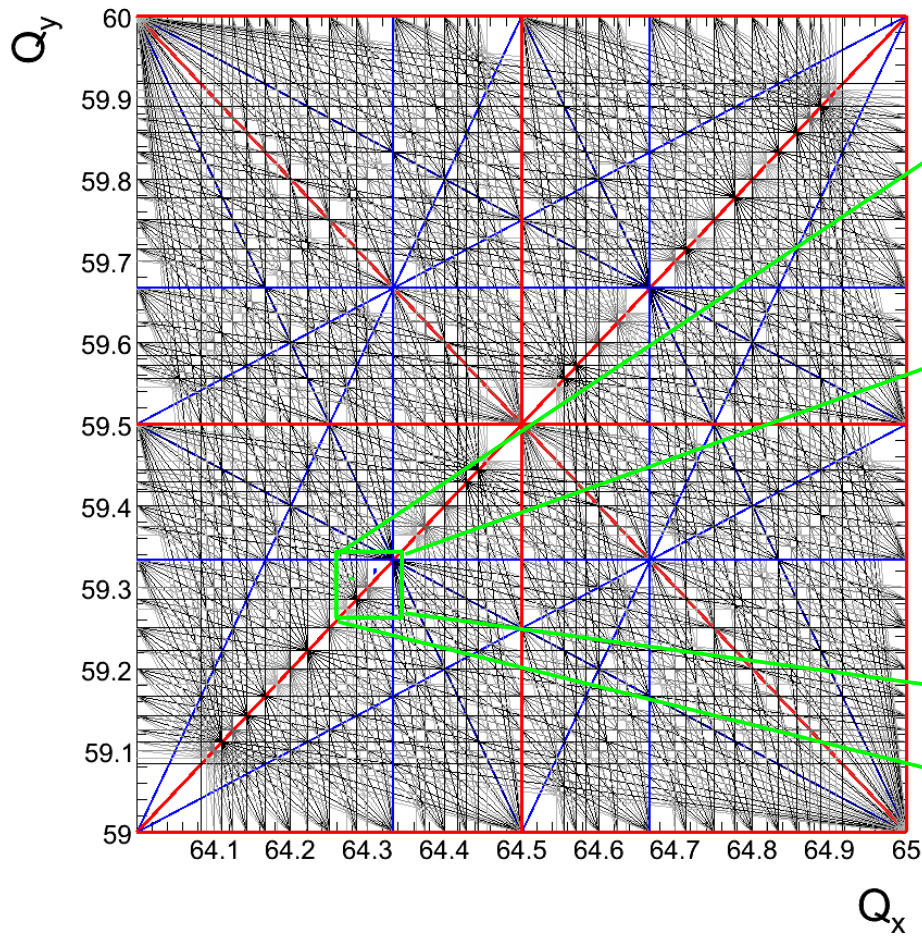


$$L = L_0 \cdot F_{crossing} \cdot \dots$$

$$F_{crossing} = \frac{1}{\sqrt{1 + \frac{\sigma_s}{\sigma_{x,y}} \tan(\theta/2)}}$$



- Example LHC:



- N.B. need to stay much further off these resonance lines due to
 - finite tune width: chromaticity, space charge, momentum spread, detuning with amplitude and resonance's stop band itself

- Small beam sizes $s(s)$ IR limited by final focus beta-function b_0 . (LHC: $b_0 = 0.55$ m)
 - max possible beta function around the detector
 - large b_{\max} : more sensitive to field errors and failures (many effects scale with β)
 - max available final focus quadrupole gradient
 - 'hour glass' effect if β^* similar to bunch length s_s :

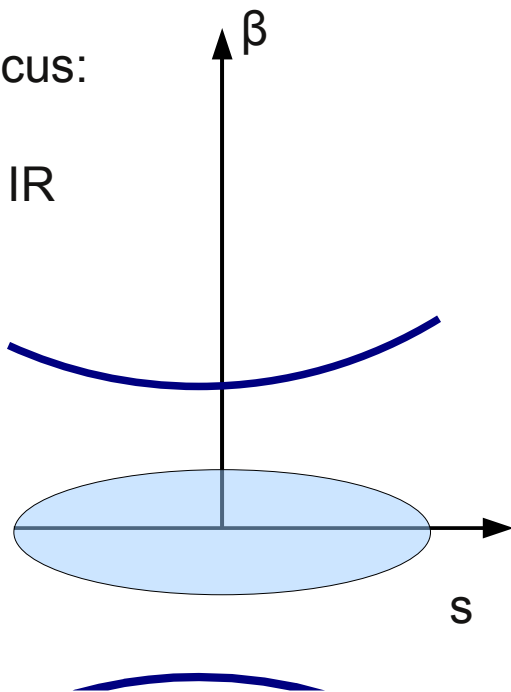
$$\sigma(s) = \sqrt{\frac{\epsilon \beta(s)}{\gamma_{rel}}}$$

$$\beta(s) = \beta_0 \left(1 + \left(\frac{s}{\beta_0} \right)^2 \right)$$

Weak final focus:

$$\beta_0 > s_s$$

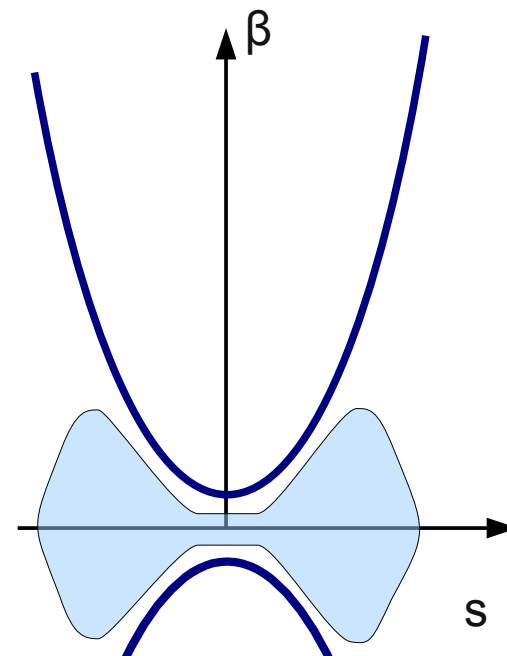
$s_s \sim \text{const @ IR}$



Strong final focus:

$$\beta_0 < s_s$$

$$s_s = s_s(s)$$



counteract with shorter bunches:

protons: decrease e_s while keeping e_{xy} constant or decreasing ... (not trivial)

- Integrated Luminosity L_{int}

$$L_{int} = \iint_0^T L(s, \epsilon, \dots, t) dt$$

- run time $t_r \approx 10$ hours (“free” parameter)
- preparation time t_p
 - magnet cycle
 - injectors,
 - detectors
 - ...
- beam lifetime τ
 - tune
 - tune spread
 - ...
 - (numerical aperture)
 - electron cloud

1st order: $\langle L \rangle \approx L_0 \cdot \tau \cdot \frac{1 - e^{-\frac{t_r}{\tau}}}{t_r + t_p}$



- Recipes to win the SUSY/Higgs Grandprix
 - optimise the machine (τ, t_p)

Picture of the Year:

